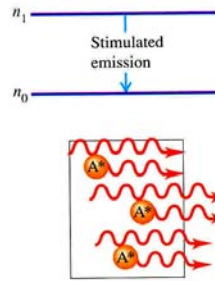


Electrons in Solids

- Today:
- Review Chapters 40, 41
 - Fermions and bosons
 - Boltzmann factor
 - Energy bands in solids
 - Fermi-Dirac distribution function
 - Some solid-state applications

Bosons

Lasers work because photons are bosons.



Atom emits a new photon into exactly the same quantum state as the original photon. This keeps happening until there is a strong beam of many photons all in the same quantum state.

Produce intense, strongly collimated laser beam.

Probability: Boltzmann Factor

Given any system in thermal equilibrium at temperature T.

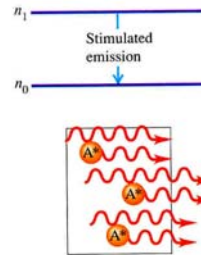
The relative probability of any two states is

$$\frac{P(E_2)}{P(E_1)} = e^{-(E_2 - E_1)/kT} = e^{-\Delta E / kT}$$

So if $\Delta E > kT$ then $P(E_2) \ll P(E_1)$

For example in a gas the number of atoms in an excited state is smaller than the number in the ground state.

Stimulated Emission



$$\Delta E = hf$$

Without pumping, at temperature T:

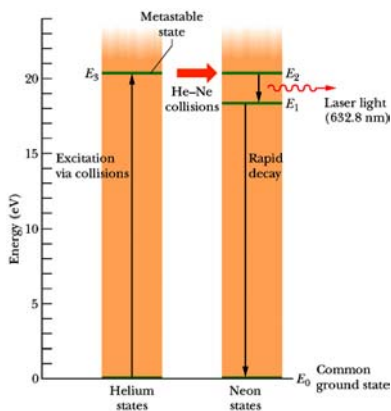
$$\frac{n_1}{n_0} = e^{-\Delta E / kT} < 1$$

Need pumping of a metastable level to produce population inversion ($n_1 > n_0$).

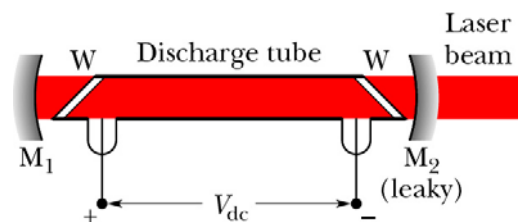
He-Ne Laser

Key to any laser is finding a metastable state which can be pumped (populated) somehow.

Get population inversion (cheat the Boltzmann factor).



Laser Construction



Fermions

- Pauli exclusion principle: No two electrons in the same state.
- That is, one electron for each set of quantum numbers: (n, l, m_l, m_s) .
- This gives the periodic table!

Q.40-3

5	6	7	8	9	10
B	C	N	O	F	Ne
10.81	12.01	14.01	16.00	19.00	20.18

Carbon ground state: $1s^2 2s^2 2p^2$

Which of the following describes the ground state of the nitrogen atom ($Z=7$)?

- (1) $1s^2 2s^2 3s^3$ (2) $1s^2 2s^2 3s^2 3p$ (3) $1s^2 2s^3 2p^2$
 (4) $1s^2 2s^2 2p^3$ (5) $1s^2 2s^2 3s 3p 3d$ (6) $1s^2 2p^5$

Q.40-3

5	6	7	8	9	10
B	C	N	O	F	Ne
10.81	12.01	14.01	16.00	19.00	20.18

Carbon ground state: $1s^2 2s^2 2p^2$

Which of the following describes the ground state of the nitrogen atom ($Z=7$)?

Begin with carbon and place one more electron in the lowest available subshell, which is 2p. Note that p orbitals can hold 6 electrons.

- (1) $1s^2 2s^2 3s^3$ (2) $1s^2 2s^2 3s^2 3p$ (3) $1s^2 2s^3 2p^2$
 (4) $1s^2 2s^2 2p^3$ (5) $1s^2 2s^2 3s 3p 3d$ (6) $1s^2 2p^5$

Ch.40 Sections 5, 6, 10

(Read these sections but don't worry about the details.)

40-5: Stern-Gerlach Experiment

- History
- Showed quantization of angular momentum

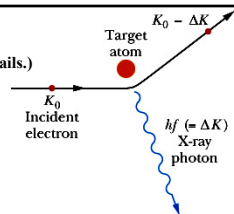
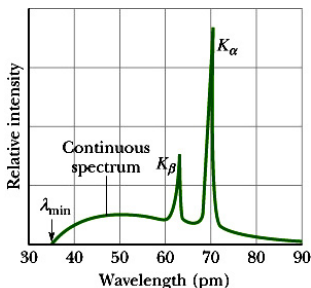
40-6: Magnetic Resonance

- Proton spins are flipped by oscillating magnetic field.
- Energy of flip is different for different molecules: NMR.
- MRI gives maps of different molecules inside the body.

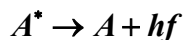
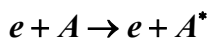
40-10: X-rays

(Read this section but don't worry about the details.)

Electron beam in an X ray tube collides with a heavy metal target, e.g. Mo.



Resulting x-ray photon energies form "continuous" and "characteristic" spectra.



X-rays (continued)

(Read this section but don't worry about the details.)

Characteristic X rays are from collisions in which a vacancy is created in an inner shell of a heavy atom.

For example in cobalt ($Z=27$) if there is a vacancy in the $n=1$ shell (K shell), then an electron from the $n=2$ shell (L shell) can jump down to fill the vacancy. The energy of the emitted photon in that case is about 7 keV (wavelength about 0.18 nm).

The jump from $n=2$ to $n=1$ is called K_α ; from $n=3$ to $n=1$ is K_β etc. Transitions down to an $n=2$ vacancy are called L_α L_β etc.

Ch. 41: Electrons in a metal



Atomic core states are full.
Conduction band is half full.



Fermi energy E_F shows highest filled state.

Easy for electron to jump to slightly higher energy state, and then move through solid.

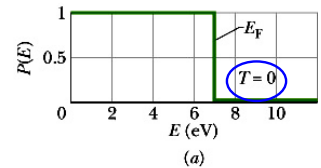
Good conductor.



Metal

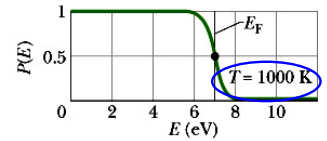
Fermi-Dirac Distribution Function

At $T=0$, states are filled just up to Fermi energy.



At $T>0$, some electrons have gotten a thermal boost.

$$P(E) = \frac{1}{e^{(E - E_F)/kT} + 1}$$



Boltzmann's constant: $k = 8.6 \times 10^{-5} \text{ eV} / \text{K}$

Example: (Based on Problem 41-5)

Consider a state in the conduction band of a metal, with energy 0.05 eV above the Fermi energy.

What is the probability that this state is occupied at room temperature?

$$kT = (8.6 \times 10^{-5} \text{ eV} / \text{K})(300 \text{ K}) = .0258 \text{ eV} \approx \frac{1}{40} \text{ eV}$$

Solution 1:

$$P(E) \approx e^{-\Delta E/kT} = e^{-\frac{.05}{1/40}} = e^{-2} = \frac{1}{e^2} = \frac{1}{7.39} = .135$$

Solution 2:

$$P(E) = \frac{1}{e^{\Delta E/kT} + 1} = \frac{1}{e^2 + 1} = \frac{1}{8.39} = .119$$

Q.41-1

Consider a metal at temperature $T = 2500 \text{ K}$. Find the probability for an electron to occupy a state in the conduction band with energy 0.20 eV above the Fermi energy: $E - E_F = 0.20 \text{ eV}$.

1. 3.5 %
2. 7 %
3. 21 %
4. 28 %
5. 35 %

Q.41-1

Consider a metal at temperature $T = 2500 \text{ K}$. Find the probability for an electron to occupy a state in the conduction band with energy 0.20 eV above the Fermi energy: $E - E_F = 0.20 \text{ eV}$.

Solution: use the Fermi-Dirac distribution function:

1. 3.5 %
 2. 7 %
 3. 21 %
 4. 28 %
 5. 35 %
- $$kT = (8.6 \times 10^{-5} \text{ eV} / \text{K})(2500 \text{ K}) = 0.215 \text{ eV}$$
- $$P(E) = \frac{1}{e^{(E - E_F)/kT} + 1} = \frac{1}{e^{0.20/0.215} + 1} = \frac{1}{2.535 + 1} = 0.28 = 28\%$$

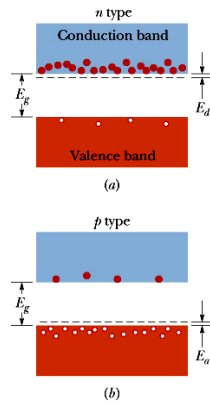
Additional topics

- Doped semiconductors
- p-n junctions
- LEDs
- Photocells

Doped semiconductors

n-type: add atom with one more valence electron. e.g. add P ($Z=15$) to Si ($Z=14$)
Extra e's lie in "donor" levels.

p-type: add atom with one less valence electron. e.g. add Al ($Z=13$) to Si ($Z=14$)
Extra empty "acceptor" levels.

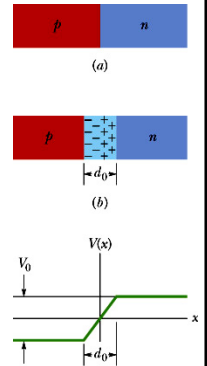


p-n Junction

Junction between p-type and n-type materials.

Diffusion across the boundary forms a depletion layer, and potential difference V_0 .

Used to form a diode. Forward voltage allows current to flow LtoR by reducing V_0 and d_0 .



LED Light-emitting diode.

Current passing through p-n junction; electrons and holes meet in the depletion layer; many recombine. This releases an energy approximately equal to the band gap.

Emission of photons!

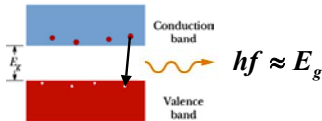
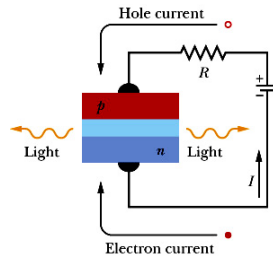
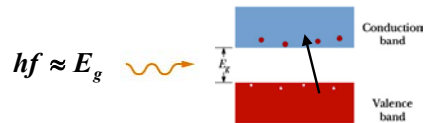


Photo-diode

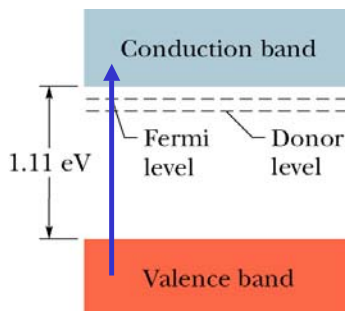
Reverse of LED. Incoming photon creates an electron-hole pair in the depletion layer of a p-n junction.

This allows flow of current; light energy is converted into electrical energy.



Silicon Photocell

Photon from sun excites electron from valence band to conduction band.



Tandem solar cell

$E_g \approx 2.2 \text{ eV}$

$E_g \approx 1.8 \text{ eV}$



Two-terminal tandem cell based on CdMnTe and HgCdTe absorbers

UT/CPEH

Two thin-film semiconductor junctions.
One band gap tuned to red light, one to blue.

Electrons in Solids

- Today:
1. Review
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 - Fermi-Dirac distribution function
 2. Some solid-state applications

Back to Solids: Chapter 41

Density of states and electrons

A crucial quantity for electronic properties of a solid is the density of states. We will not try to learn how to calculate it. Its meaning is straightforward:

$N(E) dE$ = Number of quantum states with energy in the range E to $E+dE$.

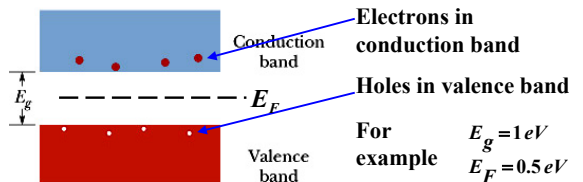


Multiply by $P(E)$ = probability of finding an electron in a state with energy E to get number of electrons with energy in range dE .

If e's were not fermions, this would be $P(E) = e^{-E/kT}$

But because e's are fermions, they obey the Fermi-Dirac probability distribution instead.

Charge carriers in semiconductor

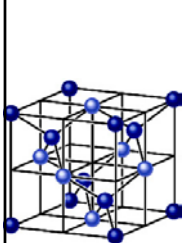


For example $E_g = 1 eV$
 $E_F = 0.5 eV$

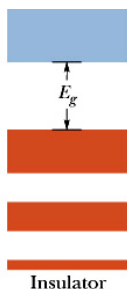
$$P(E_{conduction}) = \frac{1}{e^{(E_g - E_F)/kT} + 1} = \frac{1}{e^{(5)/.025} + 1} = \frac{1}{e^{20} + 1} = e^{-20} = 2 \times 10^{-9}$$

Multiply by Avagadro's number to get rough estimate of charge carriers per mole: $6 \times 10^{23} \times 2 \times 10^{-9} = 10^{15}$

Insulators



Diamond lattice



- All bands either full or empty.
- Band gap is large.
- Diamond:
 $E_g = 5.5 eV$

Insulator

Excellent insulator

Band gaps

Why is diamond an insulator?

Why is 5.5 eV a “large” band gap?

Electrons in filled bands cannot move because of Pauli exclusion principle.

$$E_g \gg kT$$

Must move up to conduction band.

Boltzmann factor doesn't let this happen:

Room temperature kT is about .025 eV.

$$P(E) = e^{-E/kT} = e^{-5.5/.025} = e^{-220} = 3 \times 10^{-96}$$

Also 5.5eV is larger than energy of photon of visible light. So diamond is transparent, photons are not absorbed.

Shining light on a diamond doesn't make it a conductor.