First-principles phase diagram calculations for the carbonate quasi-binary systems CaCO₃-ZnCO₃, CdCO₃-ZnCO₃, CaCO₃-CdCO₃ and MgCO₃-ZnCO₃

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Supplementary Material

| - | | | |
|----|----------------|-----------------------|-----------------------|
| | x (Å) | y (Å) | z (Å) |
| a | 4.640 | -0.008 | 4.480 |
| b | 1.897 | 4.235 | 4.480 |
| c | -0.012 | -0.008 | 6.450 |
| | \mathbf{u}_1 | u ₂ | u ₃ |
| Ca | 0.500 | 0.500 | 0.500 |
| Ca | 0.000 | 0.000 | 0.000 |
| С | 0.750 | 0.750 | 0.750 |
| С | 0.250 | 0.250 | 0.250 |
| 0 | 0.750 | 0.493 | 0.007 |
| 0 | 0.007 | 0.750 | 0.493 |
| 0 | 0.507 | 0.993 | 0.250 |
| 0 | 0.250 | 0.507 | 0.993 |
| 0 | 0.993 | 0.250 | 0.507 |
| 0 | 0.493 | 0.007 | 0.750 |

Table S1 Structural parameters of the calcite (CaCO₃) prototype used in the calculations. The other structures can be obtained through substitution and relaxation.

A. Cluster expansion formalism

In the brief description given below we provide a theoretical formalism, developed in the literature (Connolly and Williams, 1983; De Fontaine, 1994; Ducastelle, 1991; Laks et al., 1992; Sanchez et al., 1984; Zunger, 1994) over the last three decades, that allows the computation of binary and higher-order temperature-composition phase diagrams purely from first-principles calculations.

The MIT Ab-initio Phase Stability (maps) (van de Walle, 2009; van de Walle et al., 2002) code in ATAT reads in the provided lattice and automatically generates structures with different compositions at the chemically exchangeable sites. It further calculates the total energy with density functional theory (DFT) through cell volume, shape and atom position optimization and fits resulting DFT energy values to a compact polynomial representation, i.e. cluster expansion (CE),

$$E(\sigma) = \sum_{\alpha} m_{\alpha} J_{\alpha} \langle \prod_{i \in \alpha}, \sigma_i \rangle_{\alpha}$$

(Eqn. 1)

In this formula, α is a cluster of a set of exchangeable sites (*i*) in the parent lattice, and each exchangeable site is assigned a configuration variable σ_i . In a binary $M_{1-x}^A M_x^B$, or quasibinary, e.g. $M_{1-x}^A M_x^B CO_3$, σ_i takes values -1 and +1 respectively. The sum is taken over all the symmetrically inequivalent clusters of the parent lattice, and the average, denoted by angular brackets, is over all clusters symmetrically equivalent to α , m_{α} being its multiplicity. Coefficients J_{α} are called effective cluster interactions (ECIs).

The optimal cluster set and its ECIs are selected by minimizing the cross-validation (CV) score,

$$CV = \frac{1}{n} \sum_{i=1}^{n} \left(E_i - \hat{E}_{(i)} \right)^2,$$
 (Eqn. 2)

where E_i is the *ab initio* calculated energy of structure *i*, and $\hat{E}_{(i)}$ is the "leave-one-out" (without structure *i*) least-squares fitted energy to prevent over-fitting. This process is carried out iteratively when the *ab*

initio energy of a new structure is included. As the structure set grows, the cluster set typically grows as well to enhance the predictive power of the CE.

Table S2 Cluster sets and effective cluster interactions (ECIs) of the Ca_{1-x}Zn_xCO₃, Cd_{1-x}Zn_xCO₃, Ca_{1-x}Cd_xCO₃ and Mg_{1-x}Zn_xCO₃ systems.

| | | | | | ECI | ECI | ECI | ECI |
|-------|-------|----------|--------|---|---------------------|---------------------|--------------------|---------------------|
| Type | Index | Diameter | Multip | Coordinates | [Ca ₁₋ | [Cd ₁₋ | [Ca ₁₋ | [Mg ₁₋ |
| Type | muex | (Å) | licity | Coordinates | $_{x}Zn_{x}CO_{3}]$ | $_{x}Zn_{x}CO_{3}]$ | $_{x}Cd_{x}CO_{3}$ | $_{x}Zn_{x}CO_{3}]$ |
| | | | | | (meV) | (meV) |] (meV) | (meV) |
| empty | 1 | 0.000 | 1 | | 149.872 | 104.469 | 1.246 | -1.745 |
| point | 1 | 0.000 | 2 | [[0.5, 0.5, 0.5]] | 6.986 | 5.571 | -0.220 | -0.002 |
| pair | 1 | 4.097 | 6 | $\begin{bmatrix} [0.5, 0.5, 0.5], \\ [1.0, -0.0, 1.0] \end{bmatrix}$ | 7.762 | 4.001 | 1.043 | 0.900 |
| pair | 2 | 5.052 | 6 | $[[1.0, 1.0, 1.0], \\[1.0, 2.0, 0.0]]$ | -15.694 | -10.951 | -1.293 | -0.209 |
| pair | 3 | 6.450 | 6 | [[0.5, 0.5, 0.5], [0.5, 0.5, -0.5]] | -2.172 | -1.966 | 0.246 | -0.297 |
| pair | 4 | 6.505 | 6 | [[1.0, 1.0, 1.0], [1.5, 1.5, -0.5]] | 6.381 | 5.974 | -0.900 | -0.003 |
| pair | 5 | 8.193 | 6 | [[0.5, 0.5, 0.5], [-0.5, 1.5, -0.5]] | -1.832 | -1.856 | | -0.078 |
| pair | 6 | 8.236 | 12 | $\begin{bmatrix} 0.5, 0.5, 0.5 \end{bmatrix}, \\ \begin{bmatrix} -0.0, 2.0, -1.0 \end{bmatrix} \end{bmatrix}$ | -4.403 | -2.833 | | |
| pair | 7 | 8.629 | 2 | $\begin{bmatrix} [0.5, 0.5, 0.5], \\ [0.0, 0.0, 0.0] \end{bmatrix}$ | -3.071 | -1.640 | | |
| pair | 8 | 8.751 | 6 | $\begin{bmatrix} [1.0, 1.0, 1.0], \\ [-0.0, 3.0, 0.0] \end{bmatrix}$ | -1.644 | -1.285 | | |
| pair | 9 | 9.626 | 6 | [[1.0, 1.0, 1.0], [1.0, 2.0, -1.0]] | 2.370 | 1.565 | | |
| pair | 10 | 9.626 | 6 | $\begin{bmatrix} [0.5, 0.5, 0.5], \\ [0.5, 1.5, -1.5] \end{bmatrix}$ | 0.000 | 0.000 | | |
| pair | 11 | 9.999 | 12 | [[0.5, 0.5, 0.5], [-0.0, 1.0, -1.0]] | -1.452 | -1.273 | | |
| pair | 12 | 10.105 | 6 | [[1.0, 1.0, 1.0], [1.0, 3.0, -1.0]] | 0.126 | 0.365 | | |
| pair | 13 | 10.904 | 12 | [[0.5, 0.5, 0.5], [1.0, 2.0, -2.0]] | -1.108 | -0.599 | | |
| pair | 14 | 11.869 | 6 | $\begin{bmatrix} [1.0, 1.0, 1.0], \\ [1.0, 2.0, 2.0] \end{bmatrix}$ | -1.756 | -1.089 | | |

| pair | 15 | 11.988 | 6 | [[1.0, 1.0, 1.0], [-0.0, 3.0, -1.0]] | -3.290 | -2.207 | |
|------|----|--------|----|--|--------|--------|--------|
| pair | 16 | 11.988 | 6 | [[0.5, 0.5, 0.5], [-0.5, 2.5, -1.5]] | 0.000 | 0.000 | |
| trip | 1 | 5.052 | 12 | $\begin{bmatrix} [1.0, 1.0, 1.0], \\ [0.5, 1.5, 0.5], \\ [1.0, 2.0, 0.0] \end{bmatrix}$ | 0.473 | 0.695 | 0.062 |
| trip | 2 | 5.052 | 4 | [[1.0, 1.0, 1.0], [-0.0, 2.0, 1.0], [1.0, 2.0, 0.0]] | -0.207 | 0.240 | 0.000 |
| trip | 3 | 6.450 | 12 | [[0.5, 0.5, 0.5], [-0.0, 1.0, -0.0], [0.5, 0.5, -0.5]] | 0.624 | 0.237 | -0.052 |
| trip | 4 | 6.450 | 12 | $\begin{bmatrix} [0.5, 0.5, 0.5], \\ [0.5, 1.5, -0.5], \\ [0.5, 0.5, -0.5] \end{bmatrix}$ | -0.485 | -0.132 | 0.012 |
| trip | 5 | 6.505 | 12 | [[1.0, 1.0, 1.0], [0.5, 1.5, 0.5], [1.5, 1.5, -0.5]] | -1.728 | -1.359 | |
| trip | 6 | 6.505 | 12 | [[1.0, 1.0, 1.0], [1.5, 0.5, 0.5], [1.5, 1.5, -0.5]] | 0.000 | 0.000 | |
| trip | 7 | 6.505 | 12 | $\begin{bmatrix} [1.0, 1.0, 1.0], \\ [1.5, 1.5, 0.5], \\ [1.5, 1.5, -0.5] \end{bmatrix}$ | 0.717 | 0.249 | |
| trip | 8 | 8.193 | 6 | [[0.5, 0.5, 0.5], [-0.0, 1.0, -0.0], [-0.5, 1.5, -0.5]] | -3.374 | -1.169 | |
| trip | 9 | 8.193 | 12 | $\begin{bmatrix} [0.5, 0.5, 0.5], \\ [0.5, 1.5, -0.5], \\ [-0.5, 1.5, -0.5] \end{bmatrix}$ | 0.668 | -0.161 | |
| trip | 10 | 8.193 | 12 | [[0.5, 0.5, 0.5], [-0.5, 1.5, 0.5], [-0.5, 1.5, -0.5]] | 0.000 | 0.000 | |

| trip | 11 | 8.193 | 12 | [[0.5, 0.5, 0.5], [1.0, 1.0, -1.0], 0.266 [-0.5, 1.5, -0.5]] | 0.046 |
|------|----|-------|----|---|-------|
| quad | 1 | 5.052 | 4 | $\begin{bmatrix} [1.0, 1.0, 1.0], \\ [0.5, 1.5, 0.5], [-0.0, 2.0, 1.0], \\ [1.0, 2.0, 0.0] \end{bmatrix}$ | 0.996 |

B. Vibrational contributions

Vibrational free energy contributions to the phase diagrams were added through the stiffness vs length transferable force constants method (fitsvsl and svsl code in ATAT) (van de Walle and Ceder, 2002; van de Walle, 2013, 2009). It is based on the observation that bond length is typically a good predictor of bond stiffness for a given lattice and given type of bond. This approach is advantageous over the direct calculation of phonon contributions of all structures in a set because only a few structures are needed to derive the stiffness vs. length relationship. Then this relationship is used to get the (a) force constants, (b) dynamical matrices, (c) phonon spectra and (d) vibrational free energies for all structures in the CE set. Quasi-harmonic approximation considering volume expansion was used. In our calculations, we selected 10 structures to extract the information on the force constants. The minimum distance between displaced atoms was set to 9 Å, displacement of the perturbed atom to 0.2 Å, strain of the maximum volume sampled to 0.01 with 3 volume sampling points and a linear fit was performed.



Figure S1 Nearest neighbor stretching and bending stiffness constants versus bond lengths of various bonds in $Ca_{1-x}Zn_xCO_3$ and $Cd_{1-x}Zn_xCO_3$. Blue "+" and red "×" crosses stand for stretching and bending terms respectively. Lines are linear fits, which were used in the calculations of vibrational free energy contributions.

We included vibrational free energy contributions for $Ca_{1-x}Zn_xCO_3$ and $Cd_{1-x}Zn_xCO_3$, because the consolute temperature is above 1000 K, so they are likely to be significant. Figure S1 illustrates the bond stiffness vs bond length relations and the linear fits thereto. The C-O bonds in both systems are the most rigid and exhibit little scatter, but Ca-O, Cd-O and Zn-O bonds exhibit noticeable scatter. These bond stiffness values are an order of magnitude smaller than those of C-O, so they contribute to softer modes of vibration.



Figure S2 Calculated phase diagrams of (a) Ca_{1-x}Zn_xCO₃, (b) Cd_{1-x}Zn_xCO₃, (c) Ca_{1-x}Cd_xCO₃ and (d) Mg_{1-x}Zn_xCO₃. Crosses are the raw data points, and curves are interpolations and extrapolations. In (a), blue "+" and red "×" solid curves correspond to results without and with vibrational contributions, respectively. Dash-dotted curves stand for influence of the dolomite structure on the topology of the phase diagrams. Hollow circles are experimental data points. In (b), the dashed curve stands for the continuous order-disorder transition points. In (c), the dotted curve stands for the estimated consolute boundary, as demonstrated by the inset (e). In (d), the dashed curve has the same meaning as in (b). In all 4 subfigures, calcite-structure phases are labeled as I and I', and dolomite-structure phases as II. In (e), curves stand for compositions with respect to chemical potentials at various temperatures in a semi-grand-canonical ensemble simulation.

The phase diagrams are shown in Figure S2. The relevant temperatures are summarized in Table S3. For $Ca_{1-x}Zn_xCO_3$ in Figure S2(a), vibrational contributions lead to an increase in consolute temperature from 1450 K to 1600 K. The peritectoid point of the dolomite-structure phase increases from 1150 K to 1250 K. For $Cd_{1-x}Zn_xCO_3$ in Figure S2(b), the dolomite-structure phase splits through the miscibility gap between the end members, and the peritectoid point at 900 K becomes an order-disorder transition line, reaching maximum at 1050 K. Figure S3 shows the abrupt drops of the long-range order in these two systems. We conclude that the vibrational free energy contributes to the dolomite-structure phase favorably and renders it with better stability, delaying its order-disorder transition.

Table S3 Consolute temperatures (T_c), peritectoid points (T_P) and dolomite order-disorder transition temperatures (T_{trans}) of Ca_{1-x}Zn_xCO₃, Cd_{1-x}Zn_xCO₃, Cd_{1-x}Zn_xCO

| T _C without | T _C with | T _P without | T_P with | T _{trans} without | T _{trans} with |
|------------------------|---------------------|------------------------|------------|----------------------------|-------------------------|
| vib (K) | vib (K) | vib (K) | vib (K) | vib (K) | vib (K) |

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|------|---------|---------|-------|--------|------|----------|-----|---------|----------|---------|--------|---------|---------|
| | , | , | | | | , , | | , | | | | | · · · / |

| Ca _{1-x} Zn _x CO ₃ | 1450 | 1600 | 1150 | 1250 | | |
|---|-------|------|------|------|-------|------|
| $Cd_{1-x}Zn_xCO_3$ | 1000 | | 900 | | | 1050 |
| $Ca_{1-x}Cd_xCO_3$ | < 100 | | | | | |
| Mg _{1-x} Zn _x CO ₃ | | | | | < 100 | |

It is worth pointing out that because of the scatter of the bond stiffness vs bond length relations for the softer M-O bonds in Figure S1, and the closeness to 0 of dolomite structures' ΔE_f , error might have influenced the evaluation of the peritectoid point in Cd_{1-x}Zn_xCO₃. We did several trials with increasing numbers of structures to extract the information on the force constants, and large variations (200 K) occurred to the resultant consolute and transition temperatures. A high quality yet low-cost approach needs to be developed to further address this problem.



Figure S3 Long-range order as a function of temperature for Ca_{0.5}Zn_{0.5}CO₃ and Cd_{0.5}Zn_{0.5}CO₃ in canonical ensemble simulations. Blue "+" and red "×" curves correspond to results without and with the vibrational contributions, respectively. The onsets of the abrupt drops to 0.5 mark the transition temperatures of the dolomite-structure phases.

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