

**Spherical cluster method for ground state determination of site-disordered materials:
application to $\text{Ag}_x\text{Bi}_{1-x-3y}$**

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Supplemental information for V.T. Barone et. al.

Calculation of the number of combinations for a given crystal system

Let L be the number of sublattices present in the system, and let l enumerate them.

Let E_l be the number of unique elements in sublattice l , and let e_l enumerate them.

Let S_l be the total number of sites in sublattice l .

Let S_{l,e_l} be the number of sites in sublattice l of element e_l .

$$\# \text{ Combinations} = \prod_l^L \left[\prod_{e_l}^{E_l} \binom{S_l - \sum_j^{e_l-1} S_{l,j}}{S_{l,e_l}} \right]$$

The factor in parenthesis is the binomial coefficient:

$$\binom{n}{k} = \frac{n!}{k!(n-k)!}$$

It is more intuitive to follow an example.

Consider a $2 \times 2 \times 1$ cell of AgBiI_4 following the restrictions of **Table 1, row 2**. There are seven total sublattices: three of *layer 1*, three of *layer 2*, and one iodine sublattice which has no site disorder (we will therefore ignore it). Each layer contains **four** sites.

Layer 1's site distribution is **zero** Ag, **one** Bi, and **three** Va. For *Layer 1*, call the factor in square brackets N_1 .

$$N_1 = \binom{4}{0} \binom{4-(0)}{1} \binom{4-(0+1)}{3} = 4.$$

Layer 2's site distribution is **two** Ag, **one** Bi, and **one** Va. For *Layer 2*, call the factor in square brackets N_2 .

$$N_2 = \binom{4}{2} \binom{4-(2)}{1} \binom{4-(2+1)}{1} = 12.$$

Then performing the product over L (ignoring the iodine sublattice, where the factor in square brackets is 1) gives

$$\# \text{ Combinations} = N_1 N_1 N_1 N_2 N_2 N_2 = 110592$$

since there are three of each layer per cell.

Effects of Different Exchange-Correlation Functionals

Here, we list some basic properties of the minimum-energy AgBiI₄ cell when treated with a few different XC functionals. Mind, in the main text, we use the GGA-PBE functional. The calculation parameters for results displayed here follow the DFT section in the main text. We also attempted to use the MBJ functional, however, it was unable to converge to a solution even with a pre-converged wavefunction and charge density.

Functional	a (Å)	c (Å)	$\max[\mathbf{r} - \mathbf{r}_{\text{GGA}}]$ (Å)	Band Gap (eV)
GGA-PBE	4.46	21.33	0.00 (0.00)	1.71
PBEsol	4.36	20.88	0.49 (0.05)	1.67
DFT-D3	4.42	21.14	0.25 (0.11)	1.75
SCAN	4.42	21.13	0.25 (0.08)	1.90

Table S1. Properties of AgBiI₄ from **Table 1, row 1** computed with various XC functionals. a and c are lattice parameters, $\max[|\mathbf{r} - \mathbf{r}_{\text{GGA}}|]$ is the maximum deviation in the final position of any atom in a cell relaxed with some functional compared to the same atom in a cell relaxed with GGA, and the Band Gap is the minimum energy difference between the VBM (Γ for all functionals) and the CBM ((0.429, 0.000, 0.333) for all functionals). Two values are given for the entries in column four: the first is as explained above, and the second (in parentheses) removes the influence of lattice size and shape, and is therefore a measure only of atomic structure.

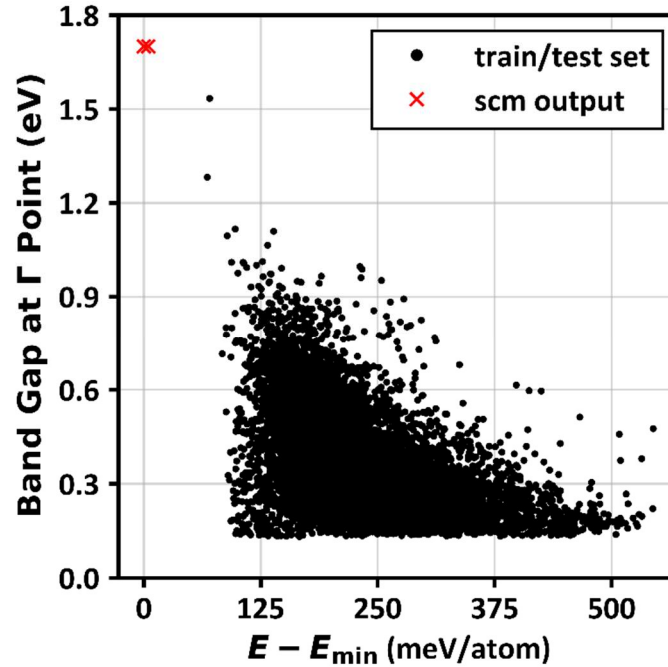


FIG. S1. Relationship between band gaps and energies of 9,202 AgBiI₄ models. Ignoring many body and relativistic effects, as well as the flat bands of silver bismuth-iodides, allow for reasonable estimates of the band gap with only the Γ point. The results suggest that very low-energy structures are necessary to reproduce the experimental (1.7 – 1.9 eV) band gap. The red crosses show the results of our method (SCM) in estimating the lowest energy structures. The two SCM outputs are nearly overlapping. The SCM was trained on 460 random data points from the 9200 points shown in the train/test set.

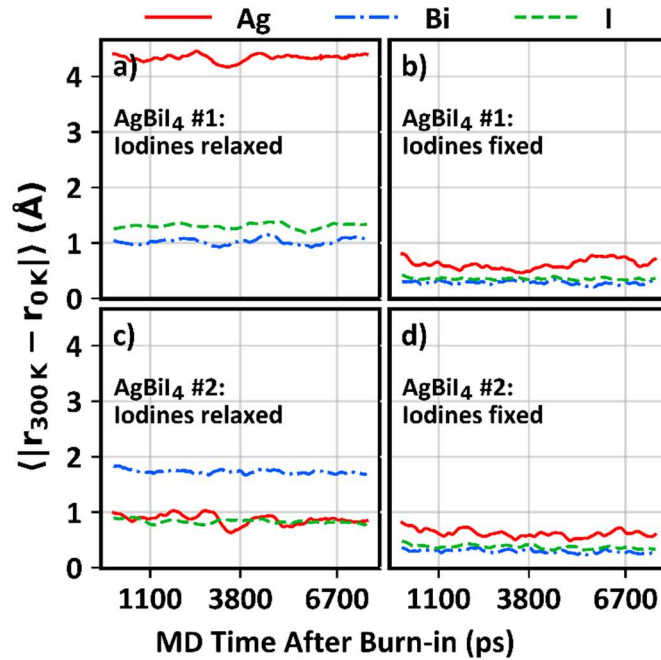


FIG. S2. Average distances between final atom positions ($r_{0\text{K}}$) and corresponding molecular dynamics (MD) positions ($r_{300\text{K}}$) for two AgBiI_4 cells. In all cases, the MD was performed at 300 K and iodine atoms were allowed to relax. **a)** and **c)** compare deviations in atom positions of the aforementioned MD with the final positions resulting from 0 K conjugate gradient (CG) minimizations. In **a)** and **c)**, the iodine atoms were allowed to relax during the CG minimizations. **b)** and **d)** are similar to **a)** and **c)** except the iodine atoms were fixed to their crystallographic positions. Comparing **a)** to **b)**, or **c)** to **d)**, shows that the CG relaxations with fixed iodine atoms more closely resemble the finite temperature MD than the CG relaxations with iodine atoms allowed to relax.

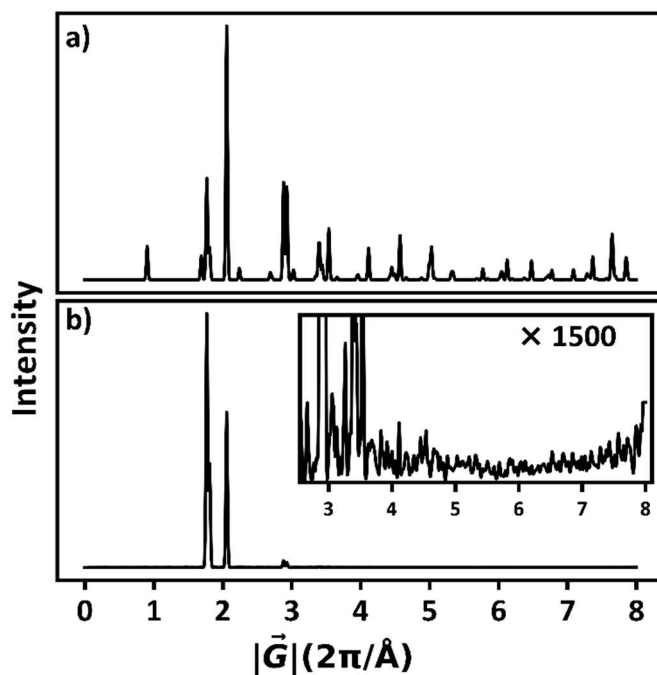


FIG. S3. a) Simulated x-ray diffraction spectra of 3000 low-energy AgBiI_4 models, keeping iodine fixed during relaxations. **b)** Simulated x-ray diffraction spectra of the same 3000 models as in **a)**, except all ions including iodine were allowed to relax. **(inset)** zoom-in of **b)** to show high-index signals (there are no high-index signals in **a)**). The y-axis is scaled by a factor of 1,500. Lorentz Polarization correction factors using $\lambda = 1.5406 \text{ \AA}$ were applied in both **a)** and **b)**. **a)** compares much more favorably to experimental data [5–9] than **b)**.

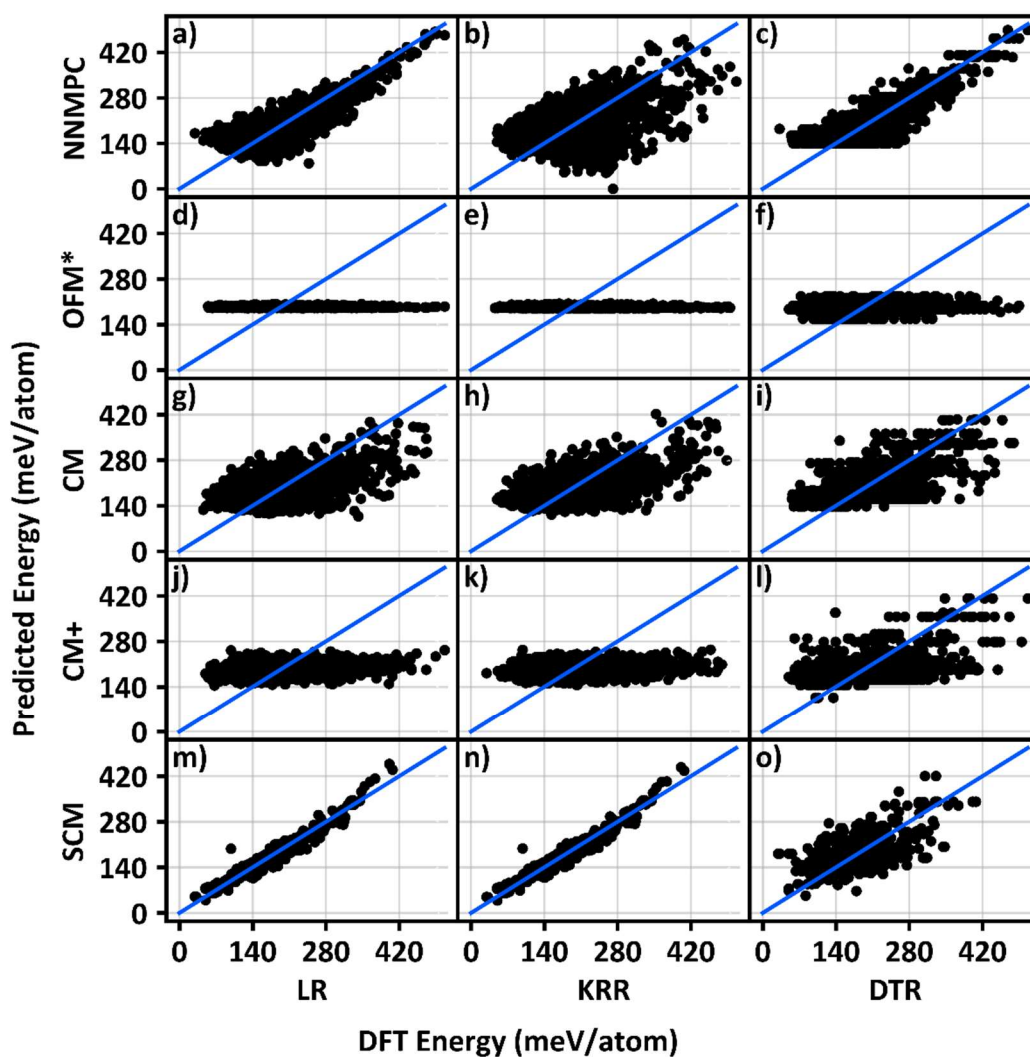


FIG. S4. Some tested energy predictions of common chemical descriptors for trigonal AgBiI_4 . In **a-l)**, an 80-20 train-test split is applied to 9200 AgBiI_4 models. In **m-o)**, the spherical cluster method (SCM) was fit to 460 points. Black data points represent testing points, and the training points are not shown. The blue line is an ideal 1:1 correlation between predicted and total energies. The descriptors are organized row-wise. From top to bottom: **a-c)** nearest-neighbor metal pair counts (NNMPC) used in our previous study of cubic AgBiI_4 [1], **d-f)** Orbital Field Matrix (OFM) eigenspectra [2], **g-i)** Coulomb Matrix (CM) eigenspectra [3], **j-l)** a periodic-cell correction to the Coulomb Matrix (CM+) eigenspectra [4], and **m-o)** the spherical cluster method (SCM)

described in this work. Three fitting techniques are shown for each descriptor, arranged column-wise: **a, d, g, j, m)** linear regression (LR), **b, e, h, k, n)** kernel ridge regression (KRR), and **c, f, i, l, o)** decision tree regression (DTR).

*: The authors of the OFM [2] suggest using the eigenspectra of the OFM created from the average of all local contributions. Here, we instead use the average eigenspectra of each local OFM to distinguish one structure from the other. The original formulation predicts the same energy for each structure.

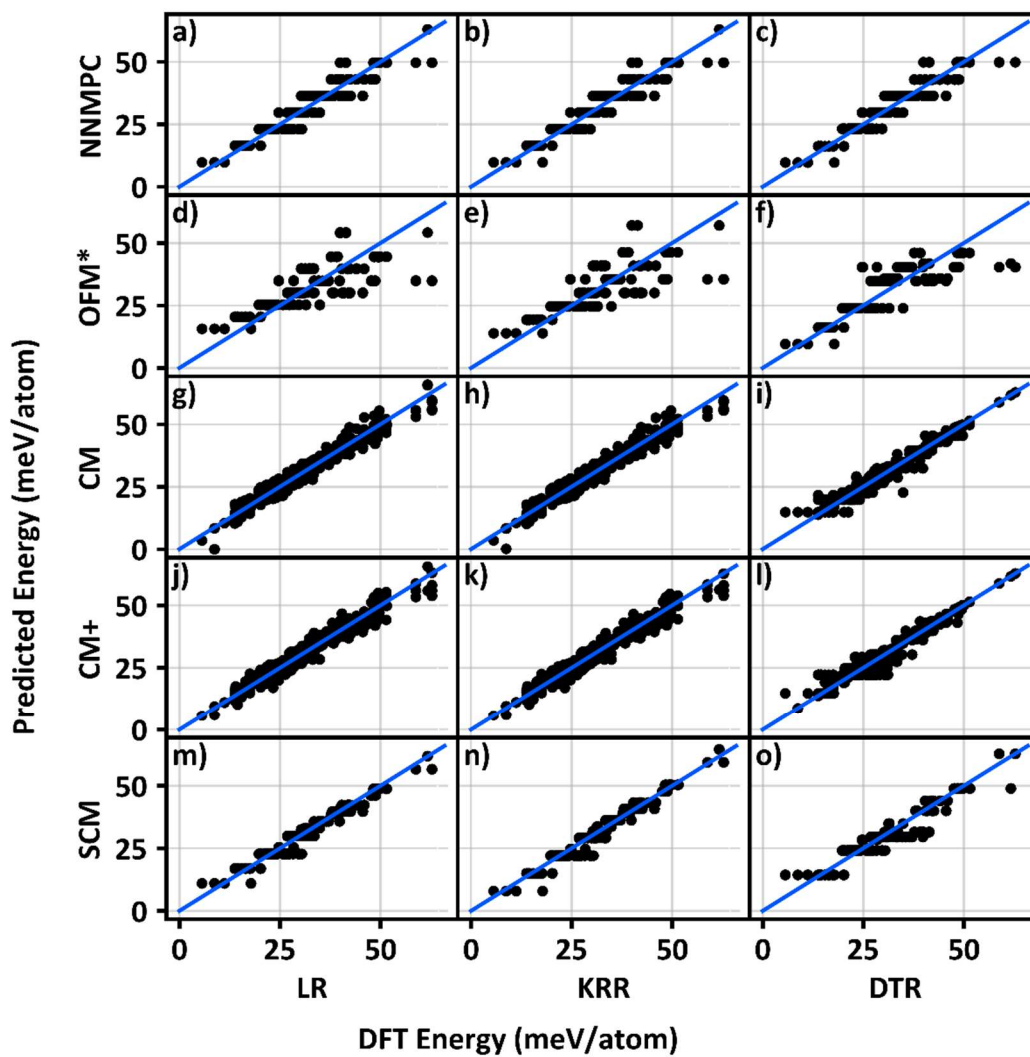


FIG. S5. Some common chemical descriptors' tested energy predictions of cubic AgBi_4 . In **a-l)**, an 80-20 train-test split is applied to all possible 12870 AgBi_4 cells described in [1]. In **m-o)**, the spherical cluster method (SCM) was fit to 30 points. Black data points represent testing points, and the training points are not shown. The blue line is an ideal 1:1 correlation between predicted and total energies. The descriptors are organized row-wise. From top to bottom: **a-c)** nearest-neighbor metal pair counts (NNMPC) used in our previous study of cubic AgBi_4 [1], **d-f)** Orbital Field Matrix (OFM) eigenspectra [2], **g-i)** Coulomb Matrix (CM) eigenspectra [3], **j-l)** a periodic-cell correction to the Coulomb Matrix (CM+) eigenspectra [4], and **m-o)** the SCM described in this work. Three fitting techniques are shown for each descriptor, arranged column-wise: **a, d, g,**

j, m) linear regression (LR), **b, e, h, k, n)** kernel ridge regression (KRR), and **c, f, i, l, o)** decision tree regression (DTR).

*: The authors of the OFM [2] suggest using the eigenspectra of the OFM created from the average of all local contributions. Here, we instead use the average eigenspectra of each local OFM to distinguish one structure from the other. The original formulation predicts the same energy for each structure.

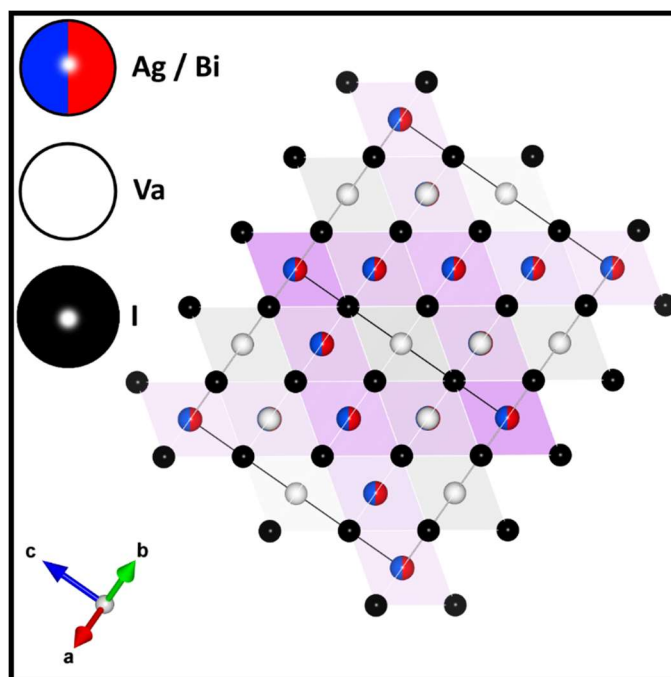


FIG. S6. The cubic unit cell of AgBiI_4 rotated to show its relation to the NaVO_2 -type models. The cell contains 32 iodine sites (black dots), 16 sites whose occupancies vary between Ag and Bi (red/blue dots), and 16 sites that are fixed as vacant (white dots).

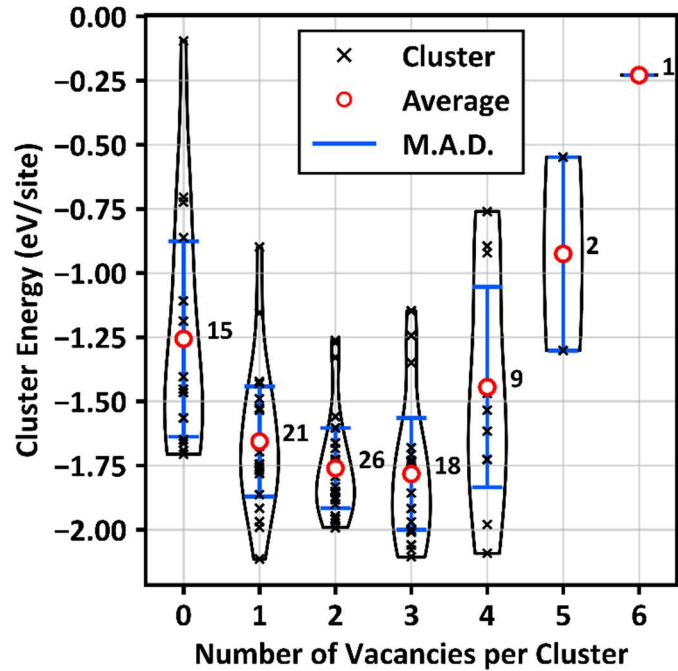


FIG. S7. Violin plot (width of black lines represent density of data points) of all 92 1st nearest neighbor (7 sites each) cluster energies vs. the number of vacant sites in each cluster. The black crosses, red circles, and blue lines represent individual data points, averages, and mean absolute deviations (M.A.D.) from the averages, respectively. The numbers to the right of each average indicate the number of data points in each group.

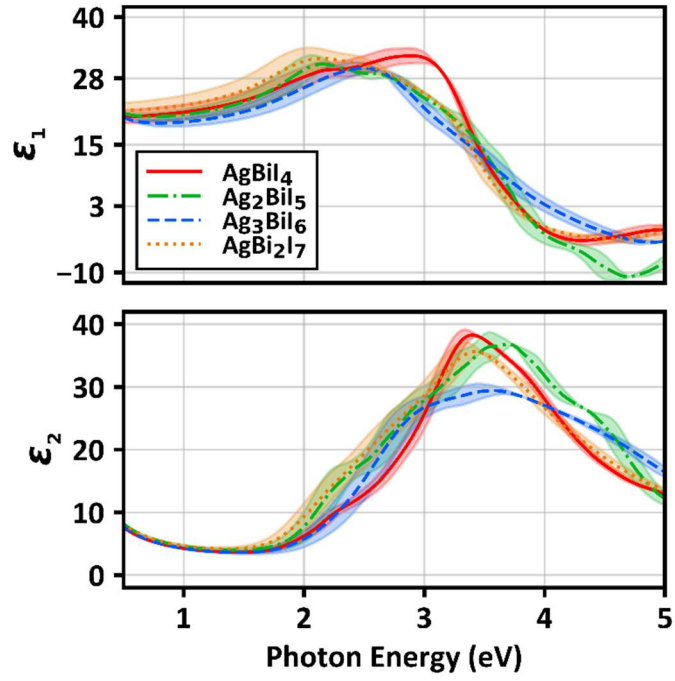


FIG. S8. PHS-corrected (PBE + HSE06 + Sum) [10] components of the complex dielectric function $\epsilon_1 + i\epsilon_2$ for the Ag-Bi-I cells at the first row of **Table I**. Filled regions represent deviations in the spectrum calculated with the PHS method for bandgap corrections of $\Delta E_g = \pm 0.1$ eV (the average error between the PBE functional and the HSE06 functional with SOC interactions included). The solid lines are averages over the filled regions.

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