Higher Order Photon Transitions in H-like and He-like Ions

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Abstract

Higher order photon transitions such as M1, M2 and two-photon decay are conveniently studied using highly-charged few electron ions. Here we discuss two examples from recent experiments which were done using the ATLAS facility at Argonne National Laboratory. The first is a test of Relativistic Quantum Mechanics involving a precision measurement of the spectral shape of the two-photon decay of the $1s2s^1S_0$ state in He-like nickel and the second is a test of the theory of damping in quantum mechanics involving observation of E1-M1 interference in the electric field quenching of metastable H-like ions.

1. Introduction

The study of atomic transition probabilities is an important part of research in atomic structure. Such studies complement measurements of transition energies because they provide more sensitivity to atomic wavefunctions and address issues such as the handling of the negative energy states [1,2,3].

Few-electron, highly charged ions provide a good testing ground for the interplay of relativistic effects and interelectron interactions. Lifetime measurements can be difficult in this regime, however, because the states are generally short lived. One solution is to study forbidden transitions from metastable states. Such states become shorter-lived as Z is increased and come into the regime of lifetimes that are convenient to study experimentally.

In this paper, we describe two recent experiments involving higher order transitions which have been done at the ATLAS heavy ion accelerator at Argonne National Laboratory. The first is a study of the spectral shape of the two-photon decay of the $1s2s^{1}S_{0}$ level in heliumlike nickel which provides a test of relativistic quantum mechanics in the intermediate Z regime. The second is a study of the angular distribution of the radiation resulting from the electric field quenching of the 2s state of H-like Ar. This experiment provides a test of the theory of damping in quantum mechanics.

2. Two-photon decay

Studies of two-photon decay are important as they test our understanding of the entire structure of an atom, since the amplitude includes contributions from all discrete and continuum states. Most of the studies of two-photon decay have relied on measurement of the lifetimes of two-photon emitting states [4,5]. For high-Z ions the lifetimes become too short to measure accurately, so we have undertaken a study of the shape of the continuum from two-photon decay. Such measurements can be done for any ion and they provide information beyond what can be obtained from lifetime

measurements which measure the transition probability integrated over the photon energies.

The spectral distribution from two-photon decay of the 1s2s 1S_0 in He-like nickel was measured in a recent experiment at ATLAS [6]. In this experiment, we were able to overcome a limitation in an earlier measurement [7] of the spectral shape in He-like Kr which was due to uncertainty in the energy-dependent efficiencies of the detectors. In the nickel experiment, two-photon decays in both H-like and He-like ions were measured in the same experiment, switching between the two species periodically. Since the continuum shape in the two-photon decay of the 2s $^2S_{1/2}$ level in H-like nickel is known precisely from theory, these data served as a calibration of the spectral efficiency of the detection system which then allowed a precision determination of the spectral distribution in the He-like two-photon transition.

In the experiment Ni^{26+} or Ni^{28+} ions were directed to the target chamber. With Ni^{26+} incident on target, the 1s2s 1S_0 level in He-like Ni was formed via excitation, whereas, with Ni^{28+} incident on the target the 2s $^2S_{1/2}$ level in H-like Ni was formed via capture. Photon coincidences associated with the decays were observed with an array of three Si(Li) X-ray detectors

A Monte Carlo code [8] was developed to simulate the actual shapes of the X-ray continua on the basis of present theories. By comparing the results of the Monte-Carlo simulation for H-like nickel with our data for the two-photon decay

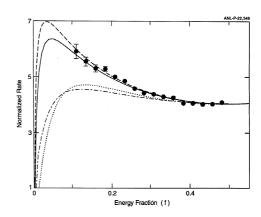


Fig. 1. Rate for two photon decay as a function of the photon energy (in units of the fraction of the transition energy). The data for He-like nickel are compared with calculations by Drake (dashed line), and Derevianko and Johnson (solid line). Also shown are Derevianko and Johnson's calculations for Z=2 (dotted line) and Z=79 (dot-dashed line). The data and theory have been divided by a factor $\omega_1(\omega_0-\omega_1)$ where ω_1 is the photon energy and ω_0 is the transition energy.

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of the $2s^2S_{1/2}$ state, we were able to constrain the simulation and obtain a model-independent determination of the shape of the continuum distribution in He-like nickel. In Fig. 1 we compare our results (points with error bars) with various calculations. The data show excellent agreement with both Derevianko and Johnson's relativistic calculations [9] and Drake's nonrelativistic calculations [10,11] for nickel. These results provide a precise measurement of the spectral distribution for two-photon decay at intermediate Z. The next step is to compare this result with a measurement in He-like gold (Z = 79) where the relativistic corrections are more important. The first data from this part of our program has already been obtained in an experiment done at GSI, Darmstadt. See the contribution by Schäffer et al. in this volume.

3. Damping Asymmetry

The metastable $2s^2S_{1/2}$ state in one-electron ions decays to the ground state either by two-photon decay or by single-photon M1 decay. If an external electric field E is applied to the ion there is, in addition, a Stark-induced E1 amplitude. Interference between the E1 and M1 amplitudes leads to an asymmetry in the angular distribution proportional to the invariant $\mathbf{k} \cdot \mathbf{E}$ which is of interest because it appears to violate time reversal (T) invariance. It has been shown however that even in a T- conserving theory, this term can be nonvanishing if damping is taken into account [12,13]. We have made the first observation of this E1 - M1damping-interference [14]. Our results provide a test of damping in quantum mechanics and can also be interpreted as a measurement of the Lamb shift in H-like Ar. E1 - M1 interference in the quenching of the $2s^2S_{1/2}$ level in polarized He⁺ ions has been measured by van Wijngaarden and Drake [15,16,17,18]. A recent review of anisotropy measurements in H-like atoms is given in Ref. [15].

In our experiment, ions formed in the metastable $2s^2S_{1/2}$ level in H-like Ar proceeded to the center of a chamber where their decay radiation was monitored by two Si(Li) X-ray detectors located on opposite sides of the beam. A magnetic field B was applied perpendicular to both the beam velocity and the detector axis. Because of their motion through this field, the ions experienced an electric field $\mathbf{E} = \gamma \frac{\mathbf{v}}{c} \times \mathbf{B}$ in their rest frame. The electric field was perpendicular to the beam velocity and parallel to the direction of the detected photons. In Fig. 2 we show typical spectra taken with one of the Si(Li) detectors at several values of magnetic field. The peak near 3 keV is a blend of Lyman- α and K α lines from H-like and He-like ions. Because of the Stark- induced E1 transitions, the H-like peak becomes more intense as the magnetic field is increased. The part of the $\mathbf{E} \cdot \mathbf{k}$ asymmetry that changes sign when the magnetic field is reversed is called r_2 . Data were taken at a number of different magnetic field strengths and fit to a straight line determining the slope:

$$r_2 = (-0.01666 \pm 0.00041) \times B(kG) \tag{1}$$

which agrees with the expected value based on the theoretical calculation[12,13]. If we interpret these data as a determination of the Lamb shift S we find $S_{exp} = 38.6(1.4) THz$ agreement with the theoretical result [19] of

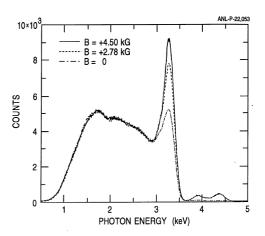


Fig. 2. Spectrum measured with a Si(Li) detector for three different magnetic field settings.

 $S_{th} = 38.2445 \, THz$. Our result is higher than but in substantial agreement with the experimental result of Gould and Marrus [20] $S_{GM} = 37.89(0.38)$ THz.

In the future, the apparatus developed for this program will be used to measure the polarization of metastable ions which have been created in an array of tilted foils [21]. Polarization would be measured by observing the angular distribution of photons arising from the quenching of the metastable ions in an electric field.

Acknowledgments

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